

Electric and Magnetic Weyl Tensors in Higher Dimensions

S. Hervik, M. Ortaggio, and L. Wylleman

Abstract Recent results on purely electric (PE) or magnetic (PM) spacetimes in n dimensions are summarized. These include: Weyl types; diagonalizability; conditions under which direct (or warped) products are PE/PM.

1 Definition and general properties

The standard decomposition of the Maxwell tensor F_{ab} into its electric and magnetic parts \mathbf{E} and \mathbf{B} with respect to (wrt) an observer (i.e., a unit time-like vector u) can be extended to any tensor in an n -dimensional spacetime [1, 2, 3]. Here we summarize the results of [3] about the Weyl tensor, and the connection with the null alignment classification [4, 5].

Consider the u -orthogonal projector $h_{ab} = g_{ab} + u_a u_b$. The “electric” and “magnetic” parts of C_{abcd} can be defined, respectively, as [3]

$$(C_+)^{ab}{}_{cd} = h^{ae} h^{bf} h_c{}^g h_d{}^h C_{efgh} + 4u^{[a} u_{[c} C^{b]e}{}_{d]f} u_e u^f, \quad (1)$$

$$(C_-)^{ab}{}_{cd} = 2h^{ae} h^{bf} C_{efk[c} u_{d]} u^k + 2u_k u^{[a} C^{b]kef} h_{ce} h_{df}. \quad (2)$$

S. Hervik
Faculty of Science and Technology, University of Stavanger, N-4036 Stavanger, Norway
e-mail: sigbjorn.hervik@uis.no

M. Ortaggio
Institute of Mathematics, Academy of Sciences of the Czech Republic
Žitná 25, 115 67 Prague 1, Czech Republic
e-mail: ortaggio@math.cas.cz

L. Wylleman
Faculty of Applied Sciences TW16, Ghent University, Galglaan 2, 9000 Gent, Belgium
e-mail: lode.wylleman@ugent.be

These extend the well-known 4D definitions [6, 7]. In any orthonormal frame adapted to u the electric [magnetic] part accounts for the Weyl components with an even [odd] number of indices u . At a spacetime point (or region) the Weyl tensor is called “purely electric [magnetic]” (from now on, PE [PM]) wrt u if $C_- = 0$ [$C_+ = 0$]. The corresponding spacetime is also called PE [PM]. Several conditions on PE/PM Weyl tensors follow.

Proposition 1 (Bel-Debever-like criteria [3]). A Weyl tensor C_{abcd} is: (i) PE wrt u iff $u_a g^{ab} C_{bc[de]u_f} = 0$; (ii) PM wrt u iff $u_{[a} C_{bc][de]u_f} = 0$.

Proposition 2 (Eigenvalues [3]). A PE [PM] Weyl operator¹ is diagonalizable, and possesses only real [purely imaginary] eigenvalues. Moreover, a PM Weyl operator has at least $\frac{(n-1)(n-4)}{2}$ zero eigenvalues.

Proposition 3 (Algebraic type [3]). A Weyl tensor which is PE/PM wrt a certain u can only be of type G , I_i , D or O . In the type I_i and D cases, the second null direction of the timelike plane spanned by u and any WAND is also a WAND (with the same multiplicity). Furthermore, a type D Weyl tensor is PE iff it is type $D(d)$, and PM iff it is type $D(abc)$.

Proposition 4 (Uniqueness of u [3]). A PE [PM] Weyl tensor is PE [PM] wrt: (i) a unique u (up to sign) in the type I_i and G cases; (ii) any u belonging to the space spanned by all double WANDs (and only wrt such u s) in the type D case (noting also that if there are more than two double WANDs the Weyl tensor is necessarily PE (type $D(d)$) [10]).

2 PE spacetimes

Proposition 5 [3]. All spacetimes admitting a shearfree, twistfree, unit timelike vector field u are PE wrt u . In coordinates such that $u = V^{-1}\partial_t$, the line-element reads

$$ds^2 = -V(t,x)^2 dt^2 + P(t,x)^2 \xi_{\alpha\beta}(x) dx^\alpha dx^\beta. \quad (3)$$

The above metrics include, in particular, direct, warped and doubly warped products with a one-dimensional timelike factor, and thus all *static* spacetimes (see also [11]). For a warped spacetime (M, g) with $M = M^{(n_1)} \times M^{(n_2)}$, one has $g = e^{2(f_1+f_2)} \left(g^{(n_1)} \oplus g^{(n_2)} \right)$, where $g^{(n_i)}$ is a metric on the factor space $M^{(n_i)}$ ($i = 1, 2$) and f_i are functions on $M^{(n_i)}$ ($M^{(n_i)}$ has dimension n_i , $n = n_1 + n_2$, and $M^{(n_1)}$ is Lorentzian).

Proposition 6 (Warps with $n_1 = 2$ [11, 3]). A (doubly) warped spacetime with $n_1 = 2$ is either type O , or type $D(d)$ and PE wrt any u living in $M^{(n_1)}$; the uplifts of the

¹ In the sense of the Weyl operator approach of [8] (see also [9]).

null directions of the tangent space to $(M^{(n_1)}, g^{(n_1)})$ are double WANDs of (M, g) . If $(M^{(n_2)}, g^{(n_2)})$ is Einstein the type specializes to $D(bd)$, and if it is of constant curvature to $D(bcd)$.

In particular, all spherically, hyperbolically or plane symmetric spacetimes belong to the latter special case.

Proposition 7 (Warps with $n_1 = 3$ [11, 3]). *A (doubly) warped spacetime with $(M^{(n_1)}, g^{(n_1)})$ Einstein and $n_1 = 3$ is of type $D(d)$ or O . The uplift of any null direction of the tangent space to $(M^{(n_1)}, g^{(n_1)})$ is a double WAND of (M, g) , which is PE wrt any u living in $M^{(n_1)}$.*

Proposition 8 (Warps with $n_1 > 3$ [11, 3]). *In a (doubly) warped spacetime*

- (i) *if $(M^{(n_1)}, g^{(n_1)})$ is an Einstein spacetime of type D , (M, g) can be only of type D (or O) and the uplift of a double WAND of $(M^{(n_1)}, g^{(n_1)})$ is a double WAND of (M, g)*
- (ii) *if $(M^{(n_1)}, g^{(n_1)})$ is of constant curvature, (M, g) is of type $D(d)$ (or O) and the uplifts of any null direction of the tangent space to $(M^{(n_1)}, g^{(n_1)})$ is a double WAND of (M, g) ; (M, g) is PE wrt any u living in $M^{(n_1)}$.*

Proposition 9 (PE direct products [3]). *A direct product spacetime $M^{(n)} = M^{(n_1)} \times M^{(n_2)}$ is PE wrt a u that lives in $M^{(n_1)}$ iff u is an eigenvector of $R_{ab}^{(n_1)}$, and $M^{(n_1)}$ is PE wrt u . (u is then also an eigenvector of the Ricci tensor R_{ab} of $M^{(n)}$, i.e., $R_{ui} = 0$.)*

A conformal transformation (e.g., to a (doubly) warped space) will not, of course, affect the above conclusions about the Weyl tensor. There exist also direct products which are PE wrt a vector u not living in $M^{(n_1)}$ [3].

Also the presence of certain (Weyl) isotropies (e.g., $SO(n-2)$ for $n > 4$) implies that the spacetime is PE, see [8, 3] for details and examples.

3 PM spacetimes

Proposition 10 (PM direct products [3]). *A direct product spacetime $M^{(n)} = M^{(n_1)} \times M^{(n_2)}$ is PM wrt a u that lives in $M^{(n_1)}$ iff all the following conditions hold (where $R_{(n_i)}$ is the Ricci scalar of $M^{(n_i)}$):*

- i) *$M^{(n_1)}$ is PM wrt u and has a Ricci tensor of the form $R_{ab}^{(n_1)} = \frac{R_{(n_1)}}{n_1} g_{ab}^{(n_1)} + u_{(a} q_{b)}$ (with $u^a q_a = 0$)*
- ii) *$M^{(n_2)}$ is of constant curvature and $n_2(n_2 - 1)R_{(n_1)} + n_1(n_1 - 1)R_{(n_2)} = 0$.*

Further, $M^{(n)}$ is PM Einstein iff $M^{(n_1)}$ is PM Ricci-flat and $M^{(n_2)}$ is flat.

See [3] for explicit (non-Einstein) examples. However, in general PM spacetimes are most elusive. For example,

Proposition 11 (3). *PM Einstein spacetimes of type D do not exist.*

In [3] also several results for PE/PM Ricci and Riemann tensors have been worked out, along with corresponding examples. In general, we observe that PE/PM tensors provide examples of *minimal tensors* [12]. Thanks to the *alignment theorem* [13], the latter are of special interest since they are precisely the *tensors characterized by their invariants* [13] (cf. also [3]). This in turn sheds new light on the classification of the Weyl tensor [5], providing a further invariant characterization that distinguishes the (minimal) types G/ID from the (non-minimal) types II/III/N.

Acknowledgments

M.O. acknowledges support from research plan RVO: 67985840 and research grant no P203/10/0749.

References

1. J.M.M. Senovilla, *Super-energy tensors*, Class. Quantum Grav. **17**, 2799 (2000)
2. J.M.M. Senovilla, *General electric-magnetic decomposition of fields, positivity and Rainich-like conditions*, in *Reference Frames and Gravitomagnetism*, ed. by J.F. Pascual-Sánchez, L. Floría, A. San Miguel, F. Vicente (World Scientific, Singapore, 2001), pp. 145–164
3. S. Hervik, M. Ortaggio, L. Wylleman, *Minimal tensors and purely electric or magnetic spacetimes of arbitrary dimension*, ArXiv e-prints [arXiv:1203.3563 [gr-qc]] (2012)
4. R. Milson, A. Coley, V. Pravda, A. Pravdová, *Alignment and algebraically special tensors in Lorentzian geometry*, Int. J. Geom. Meth. Mod. Phys. **2**, 41 (2005)
5. A. Coley, R. Milson, V. Pravda, A. Pravdová, *Classification of the Weyl tensor in higher dimensions*, Class. Quantum Grav. **21**, L35 (2004)
6. A. Matte, *Sur de nouvelles solutions oscillatoires de équations de la gravitation*, Can. J. Math. **5**, 1 (1953)
7. H. Stephani, D. Kramer, M. MacCallum, C. Hoenselaers, E. Herlt, *Exact Solutions of Einstein's Field Equations*, 2nd edn. Cambridge Monographs on Mathematical Physics (Cambridge University Press, Cambridge, 2003)
8. A. Coley, S. Hervik, *Higher dimensional bivectors and classification of the Weyl operator*, Class. Quantum Grav. **27**, 015002 (2010)
9. A. Coley, S. Hervik, M. Ortaggio, L. Wylleman, *Refinements of the Weyl tensor classification in five dimensions*, Class. Quantum Grav. **29**, 155016 (2012)
10. L. Wylleman, *On Weyl type II or more special spacetimes in higher dimensions*. In preparation
11. V. Pravda, A. Pravdová, M. Ortaggio, *Type D Einstein spacetimes in higher dimensions*, Class. Quantum Grav. **24**, 4407 (2007)
12. R.W. Richardson, P.J. Slodowy, *Minimum Vectors for real reductive algebraic groups*, J. London Math. Soc. **42**, 409 (1990)
13. S. Hervik, *A spacetime not characterized by its invariants is of aligned type II*, Class. Quantum Grav. **28**, 215009 (2011)